We have computed Mössbauer spectra from two Lorentz lines. The isomer shift is shown in figure 2. We suggest that line 1 corresponds to the crystal phases Mg<sub>2</sub>Sn I (0-35 kbar) and II (>35 kbar), and the steep change of the isomer shift  $\delta E_{tr}$  provides information on the change of the electronic configuration of Mg<sub>2</sub>Sn at the phase transition. The change in the isomer shift  $\delta E_{tr}$  reflects the increase of electron density at Sn nuclei, which is a usual feature of a first-order phase transition. The decrease of  $\delta E$  preceding the phase transition is also a common characteristic for all the tin systems studied so far and shows a decrease of electron density at Sn nuclei before transition.

It is too early to reach a definite conclusion about the nature of line 2 which occurs in the transition pressure region, but it is possible that this line is related to some intermediate state of the  $Mg_2Sn$  system present at the grain boundaries of  $Mg_2Sn$  I and  $Mg_2Sn$  II phases. It should be remembered that the Mössbauer effect can show up both crystalline and amorphous systems.

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Investigation of S-state splitting the gadolinium ion in fluorite crystals by ESR and ENDOR at high pressure

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Splitting of the S ground state (with orbital angular momentum L = 0) of paramagnetic ions has been investigated by many workers since the appearance of the first paper on the subject by Van Vleck and Penny (1934). Watanabe (1957) specifies twelve processes which can lead to the splitting of the S state. Powell *et al.* (1960) and Gabriell *et al.* (1961) consider both odd and even perturbation processes by the crystalline field potential Dq. The idea of identifying actually occurring processes with the aid of experiments under high hydrostatic pressures was put forward by Watanabe (1960).

We have studied the influence of high pressure on the electron spin resonance (ESR) spectra of the gadolinium ion in fluorite crystals. The logarithmic derivative, n, of the cubic splitting parameter  $b_4$  with respect to the lattice parameter a is equal to  $7 \cdot 0 \pm 0.3$  for CaF<sub>2</sub> (Kasatochkin *et al.*, 1973). The quantity n can be regarded as the exponent in the relationship

$$b_4 \propto a^{-n} \tag{1}$$

approximating the behaviour of  $b_4$  as a function of a.

The value of the exponent n is of use only when the local compressibility of the paramagnetic complex is known. In order to determine the local compressibility we have made high-pressure determinations by electron-nuclear double resonance (ENDOR) on fluorine ligands (Kasatochkin and Yakovlev, 1975). The logarithmic derivative m of the anisotropic hyperfine coupling parameter  $A_p$  with ligands of the first coordination shell with respect to the lattice parameter a is equal to  $2 \cdot 26 \pm 0 \cdot 1$ . The quantity m has a similar meaning to the quantity n in equation (1). The difference between the value of m and the dipole-dipole value m = 3 is related to the difference between local and bulk compressibilities, if we neglect the small contribution of anisotropic exchange.

The value of *n* corrected for the difference between local and bulk compressibilities is 9.3. This value is reasonably close to the value of 10 yielded by the point-charge model for the case when the cubic splitting parameter  $b_4$  is proportional to the square of the crystalline field parameter Dq:

 $b_4 \propto (Dq)^2$ .

Thus, using the  $CaF_2$ :  $Gd^{3^+}$  system as an example, we have solved the problem posed by Watanabe (1960).

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## High pressure polymorphs of niobium dichalcogenides

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Specimens of niobium disulphide, niobium diselenide, and niobium ditelluride were synthesised from stoichiometric mixtures of the components at 90 kbar (table 1). At 1200-1350°C two new phases were obtained: NbS<sub>2</sub> II and NbSe<sub>2</sub> II. All lines on the powder photographs can be assigned on the basis of a hexagonal unit cell with parameters a = 3.33 Å, c = 23.8 Å,  $c/a = 4 \times 1.787$  (NbS<sub>2</sub> II); and a = 3.46 Å, c = 24.8 Å,  $c/a = 4 \times 1.791$  (NbSe<sub>2</sub> II). These phases were found to be isostructural with 4H(b)-TaSe<sub>2</sub> (Brown and Beerntsen, 1965) with the same

Temperature (°C)	Starting materials				
	Nb+2S	Nb+2Se	Nb+2Te		
<600	no reaction	no reaction	no reaction		
800-1000	$3R-NbS_2$ $3R-NbS_2 + 2H-NbS_2$	NbSe <sup>a</sup> <sub>4</sub> + Nb	NbTe <sup>b</sup> <sub>4</sub> + Nb		
1200-1350	$\begin{array}{cccc} 00-1350 & NbS_2 \ II & NbSe_2 \ II & NbSe_4 + Nb \end{array}$		$NbTe_4 + Nb_5Te_4$		

Table 1. Synthesis of niobium dichalcogenides at 90 kbar.

<sup>a</sup>  $\beta$ -NbSe<sub>4</sub> (Selte and Kjekshus, 1965).

<sup>b</sup> High pressure phase NbTe<sub>4</sub> II (Bjerkelund et al., 1968).

Table 2. Interatomic distances (Å) for the new polymorphs of niobium dichalcogenides.

	Nb-X (trigonal-prismatic coordination)	Nb-X (octahedral coordination)	X-X (between layers)	
NbS <sub>2</sub> II	2.48	2.47	3.42	
NbSe <sub>2</sub> II	2.58	2.57	3.56	